

# Formation and recombination of a fully ionized plasma by nanosecond discharges in ambient conditions

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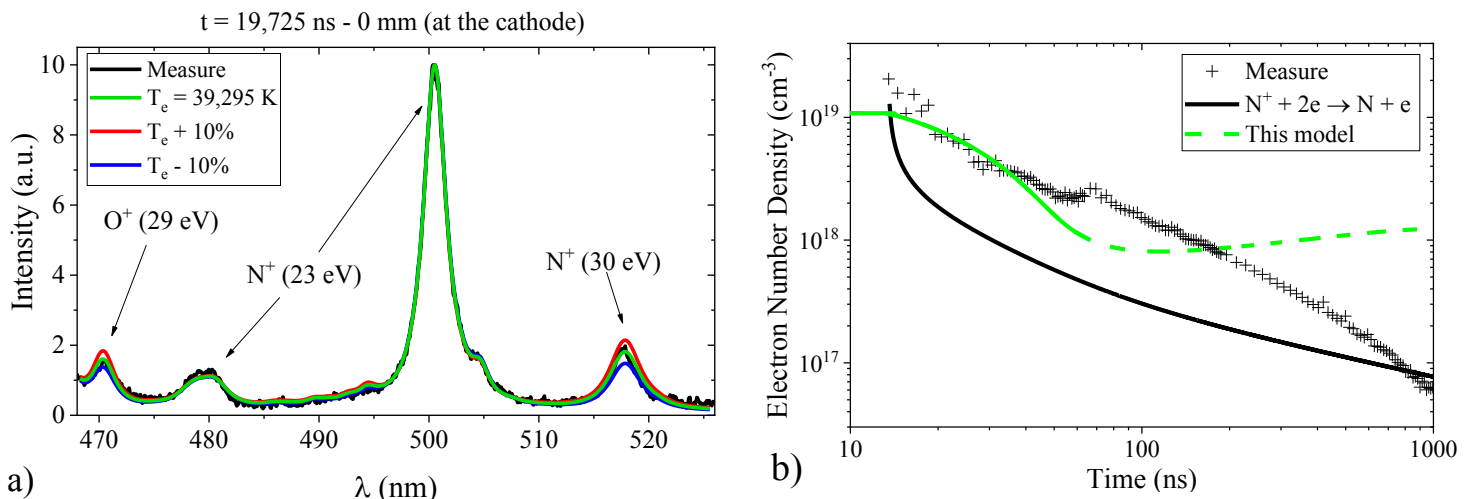
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## Résumé

Nanosecond Repetitive Pulse (NRP) discharges in pin-to-pin configurations are classified in three regimes: corona, glow and spark [1]. These three regimes are weakly ionized ( $< 1\%$ ) and out of equilibrium. It was however shown [2],[3] that nanosecond discharges can be fully ionized. The formation of the corona, glow and spark regimes are well known [1], but the formation of a fully ionized plasma by a nanosecond pulse is still partly understood. In this presentation, we will describe the steps of formation of the fully ionized discharge in a 2-mm gap and the recombination of the electrons. The imaging of the discharge was performed from 50 to 1000 mbar with a sub-nanosecond iCCD gates. The following steps of the thermal spark formation has been seen (time stamps in brackets are representative of atmospheric pressure):

1. A homogeneous plasma channel is first formed between the electrodes.
2. A filament is formed at the cathode (5.25 ns), and then at the anode (6.25 ns).
3. The filaments propagate toward each other. Typical speed is  $2 \times 10^4$  m/s at the anode and  $4 \times 10^5$  m/s at the cathode.
4. At 10.25 ns, a fully ionized channel connects the two electrodes.

Optical Emission Spectroscopy of the discharge was performed and we used Stark broadening of the  $N^+$ ,  $O^+$  and  $H_\alpha$  lines for electron number density measurements. A typical fit of the  $N^+$  and  $O^+$  lines is shown in Fig 1. a). These measurements were spatially resolved in the interelectrode gap. We show that: (i) the filaments are fully ionized and (ii) the electron number density is homogeneous along the filaments. In consequence, the plasma is homogeneously fully-ionized by the end of the pulse. These measurements are extended in Fig 1 b) with Stark broadening of  $H_\alpha$ . The excitation electron temperature,  $T_e$ , was also measured using the relative ratio of the observed lines of  $N^+$ . We artificially added 10% to the electronic temperature of the best fits in Fig. 1 b) to estimate the relative error in the measurement of  $T_e$ . Using these temperatures, we could run a 0D model of the afterglow kinetics. We show electron- $N^+$  recombination of electrons is responsible for the electron number density decay. However, ionization of  $N(2D)$  by electrons has to be taken into account to reproduce the experimental results. It is shown in Fig. 1 b) that better agreement is achieved in the first 40 ns with a model including ionization than with only 3-body recombination.



**Figure 1. a)** Typical fit (in green) of the  $N^+$  and  $O^+$  emission (black). Continuum emission was subtracted (assumed to be linear). We obtained  $n_e = 1.4 \times 10^{19} \text{ cm}^{-3}$  and  $T_e = 39.2895 \text{ K}$  with the best fit. Two other simulations (in red and blue) show the error in  $T_e$  is less than 10%. **b)** Measured decay of the electron number density (crosses) compared with electron-ion recombination only (black solid line) and the 0D model of this abstract (green).

## Références

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- [2] N. Minesi, S. A. Stepanyan, P. B. Mariotto, G.-D. Stancu, and C. O. Laux, *AIAA Scitech 2019 Forum*, 2019.
- [3] A. Lo et al., *Plasma Sources Sci. Technol.*, vol. 26, no. 4, p. 45012, 2017.